

# QGP BUBBLE DYNAMICS IN THE PRESENCE OF MAGNETIC FIELD

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Emergent Topics in relativistic Hydrodynamics,  
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S. Namitha, Lakshmi J. Naik and V. Sreekanth,  
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*manuscript under preparation*

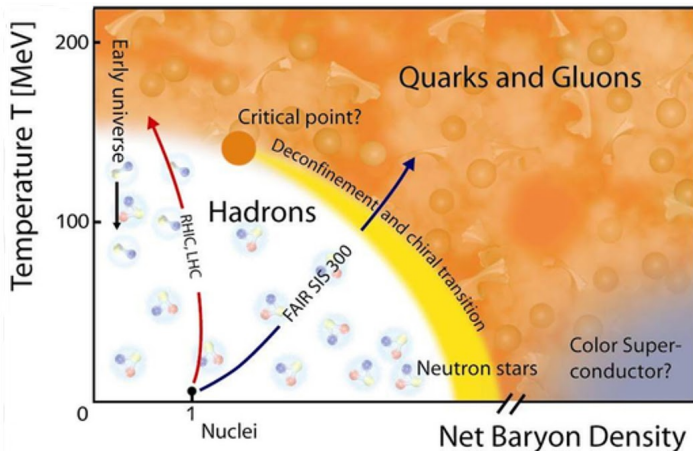


Figure: QCD phase diagram

- Nuclear collisions to be performed at FAIR and NICA are expected to create QGP at low temperature and high baryon densities
- In such collisions, the quark-hadron conversion is a first order phase transition - nucleation theory can be employed
- Hadronic matter requires some time to convert to QGP
- Compressed hadronic matter  $\longrightarrow$  QGP bubbles

# Hadron $\longrightarrow$ QGP phase transition

- The transition involves formation of QGP bubbles in hadronic medium followed by collapse of Hadron gas (HG) bubbles in QGP.

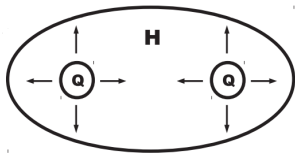


Figure: Expansion of QGP bubbles.

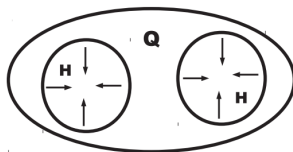
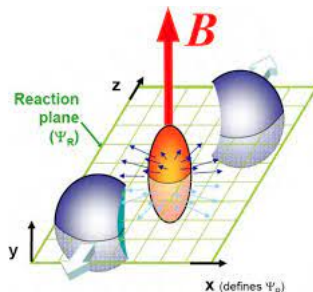
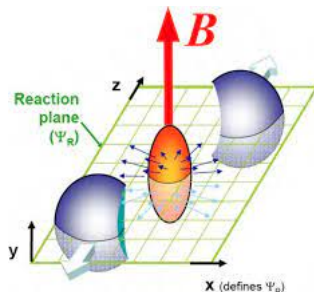


Figure: Collapse of HG bubbles.

Figs. from D. A. Fogaca et al., PRC 93, 055204 (2016)



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- Matter created in heavy ion collisions is subjected to large magnetic field produced by the spectators ( $\sim eB = 10^{19}$  G)
- **AIM:** Effect of magnetic field on expansion or collapse of QGP/HG bubbles in a relativistic treatment.

# Rayleigh-Plesset Equation

- Collapse and expansion of bubbles in an external medium can be modelled using the Rayleigh-Plesset (RP) Equation
- Collapse of an empty bubble in liquid : Rayleigh equation

$$R\ddot{R} + \frac{3}{2}\dot{R}^2 = \frac{p_i - p_\infty}{\rho}$$

$p_\infty$  : far-field pressure [Lord Rayleigh, Philos. Mag. 34, 94 (1917)]

- Non-relativistic, incompressible version of RP equation :  
[M. S. Plesset, (1949)]

$$R\ddot{R} + \frac{3}{2}\dot{R}^2 = \frac{1}{\rho} \left( p_i - p_0 - \frac{2\sigma}{R} - \frac{4\eta\dot{R}}{R} \right)$$

$p_0$  - static pressure outside the bubble wall

$p_i$  - constant pressure inside the bubble

$\sigma, \eta$  - surface tension & shear viscosity



# Relativistic RP Equation

- Relativistic version of the RP equation (RRP) was done by [Elze \*et al\*](#)
- Equation was extended to the relativistic regime at finite temperature by [D. A. Fogaca \*et al.\*](#)
- Lagrangian for a spherically symmetric system

$$\mathcal{L} = -4\pi \int_0^\infty dr r^2 \epsilon(n); \text{ where } n = \rho/\gamma.$$

- Contributions of internal and external matter :

$$\begin{aligned}\mathcal{L} &= \mathcal{L}_{in} + \mathcal{L}_{ex} \\ &= -4\pi \int_0^R dr r^2 \epsilon_{in}(n_{in}) - 4\pi \left\{ \lim_{R_\infty \rightarrow \infty} \int_R^{R_\infty} dr r^2 \epsilon_{ex}(n_{ex}) \right\}\end{aligned}$$

[Elze et al., J. Phys. G 25, 1935 \(1999\)](#) [D. A. Fogaca et al., PRC 93, 055204 \(2016\)](#)

- Number density profile for each phase :

$$\rho_{in} = \frac{3N_{in}}{4\pi R^3}, \quad 0 < r < R$$

$$\rho_{ex} = \frac{3N_{ex}}{4\pi(R_{\infty}^3 - R^3)}, \quad R < r < R_{\infty}$$

- Now the velocity profiles can be obtained as

$$v_{in} = \frac{r}{R} \dot{R}, \quad 0 < r < R$$

$$v_{ex} = \frac{(R_{\infty}^3 - r^3)}{(R_{\infty}^3 - R^3)} \frac{R^2}{r^2} \dot{R}, \quad R < r < R_{\infty}$$

- Temperature is introduced with the thermodynamic relation :

$$d\epsilon = T ds + \mu dn; \quad \mu = \frac{\partial \epsilon}{\partial n}$$

- Euler-Lagrange equation gives the relativistic RP equation :

$$\frac{d}{dt} \left[ (I_1 + I_2) R^3 \dot{R} \right] + (I_1 - 2I_2) R^2 \dot{R}^2 = (P_{in}^* - P_{ex}^*) R^2$$

$$I_1 = \int_0^1 dx x^4 (\epsilon_{in} + P_{in}^*) \gamma_{in}^2$$

$$P_{in}^* = P_{in} - \frac{\partial P}{\partial T} T|_{in}$$

$$I_2 = \int_1^\infty \frac{dx}{x^2} (\epsilon_{ex} + P_{ex}^*) \gamma_{ex}^2$$

$$P_{ex}^* = P_{ex} - \frac{\partial P}{\partial T} T|_{ex}$$

$$x = r/R.$$

D. A. Fogaca et al., PRC 93, 055204 (2016)

- Matter created in heavy ion collisions is subjected to large magnetic field produced by the spectators ( $\sim eB = 10^{19}$  G)
- Approximations : External magnetic field is constant and homogeneous
- Medium surrounding the bubble is in-compressible, irrotational
- Effects due to magnetic friction is included through dissipation function

$$\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{R}} \right) = \frac{\partial \mathcal{L}}{\partial R} - \frac{\partial \mathcal{F}}{\partial \dot{R}}$$

- Current induced in the medium outside the bubble :

$$\mathbf{j} = \sigma[(\mathbf{v} \times \mathbf{B}) + \mathbf{E}] = \sigma[(\mathbf{v} \times \mathbf{B}) - \nabla\psi]$$

$$[\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} = 0]$$

- Magnetic force due to the current

$$\begin{aligned}\mathbf{f} = \mathbf{j} \times \mathbf{B} &= \sigma[(\mathbf{v} \times \mathbf{B}) \times \mathbf{B}] \\ &= -\sigma[\mathbf{v}\mathbf{B}^2 - \mathbf{B}(\mathbf{v} \cdot \mathbf{B})]\end{aligned}$$

B.O. Kerbikov et. al PRD101,094004 (2020)

- Energy dissipation due to magnetic friction :

$$\mathcal{Q} = -\mathbf{f} \cdot \mathbf{v} = 4\pi\sigma \mathbf{B}^2 \int_R^{R_\infty} v_{\text{ex}}^2 r^2 dr = 4\pi\sigma \mathbf{B}^2 R^3 \dot{R}^2 = 2\mathcal{F}$$

- Using the dissipation function obtained above, relativistic RP with magnetic field is obtained as

$$\frac{d}{dt} \left[ (I_1 + I_2) R^3 \dot{R} \right] + (I_1 - 2I_2) R^2 \dot{R}^2 = (P_{\text{in}}^* - P_{\text{ex}}^*) R^2 + \sigma B^2 R^3 \dot{R}$$

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- Using the dissipation function obtained above, relativistic RP with magnetic field is obtained as

$$\frac{d}{dt} \left[ (l_1 + l_2) R^3 \dot{R} \right] + (l_1 - 2l_2) R^2 \dot{R}^2 = (P_{in}^* - P_{ex}^*) R^2 + \sigma B^2 R^3 \dot{R}$$

- RRP equation is solved employing the lattice QCD EoS for QGP phase and the Ideal HRG for HG matter.

$$P(T, \mu) = T^4 \int_0^T dT' \frac{e^{-h_1/\tau' - h_2/\tau'^2}}{T'} \left[ h_0 + \frac{f_0 [\tanh(f_1 \tau' + f_2) + 1]}{1 + g_1 \tau' + g_2 \tau'^2} \right] + \frac{\chi_2}{2} \mu^2 T^2$$

$$\epsilon(T, \mu) = 3P(T, \mu) + \frac{\mu^2}{2} T^3 \frac{d\chi_2}{dT} + T^4 e^{-h_1/\tau - h_2/\tau^2} \left[ h_0 + \frac{f_0 [\tanh(f_1 \tau + f_2) + 1]}{1 + g_1 \tau + g_2 \tau^2} \right]$$

where,

$$\chi_2(T) = e^{-h_3/\tau - h_4/\tau^2} f_3 [\tanh(f_4 \tau + f_5) + 1]$$

$\tau = T/T_c$  with  $T_c = 200$  MeV and  $\mu = 350$  MeV.

[S. Borsanyi et al., PLB 730, 99 (2014)]



Pressure and energy density of Ideal HRG :

$$\begin{aligned}P(T, \mu) &= \sum_i p_i(T, \mu_i) \\&= \sum_i \frac{d_i}{6\pi^2} \int dm f_i(m) \int_0^\infty \frac{dk k^4}{\sqrt{k^2 + m^2}} \\&\quad \times \left[ \exp\left(\frac{\sqrt{k^2 + m^2} - \mu_i}{T}\right) + \eta_i \right]^{-1} \\ \epsilon(T, \mu) &= \sum_i \epsilon_i(T, \mu_i) \\&= \sum_i \frac{d_i}{2\pi^2} \int dm f_i(m) \int_0^\infty k^2 dk \sqrt{k^2 + m^2} \\&\quad \times \left[ \exp\left(\frac{\sqrt{k^2 + m^2} - \mu_i}{T}\right) + \eta_i \right]^{-1}\end{aligned}$$

(1)

- Summation over all strange and non-strange hadrons listed in Particle Data Tables (mesons upto  $f_2(2340)$  and baryons upto  $N(2600)$ ) [K. A. Olive et al, Chin. Phys. C 38, 090001 (2014)]
- $f_i(m) = \delta(m - m_i)$ , where  $m_i$  is the mass of stable hadrons

- RRP equation is solved numerically employing the lattice QCD EoS for QGP phase and the Ideal HRG for HG matter

$$\frac{d}{dt} \left[ (l_1 + l_2) R^3 \dot{R} \right] + (l_1 - 2l_2) R^2 \dot{R}^2 = (P_{in}^* - P_{ex}^*) R^2 + \sigma B^2 R^3 \dot{R}$$

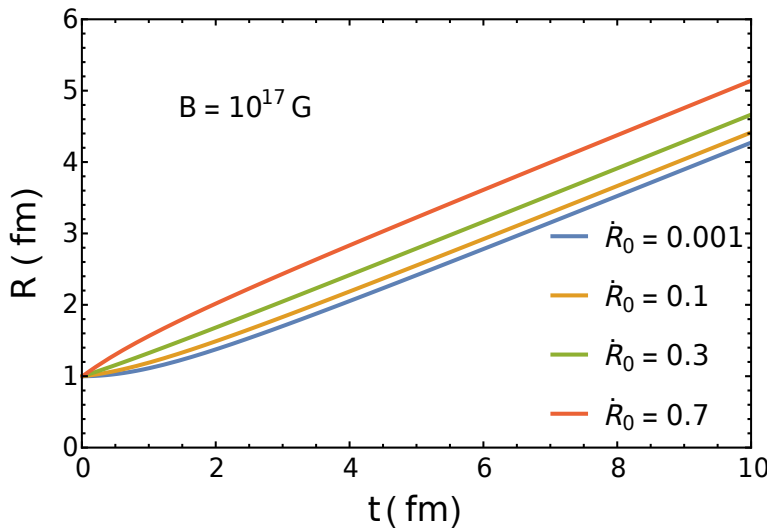


Figure: Expansion of QGP bubble in hadronic medium for  $T_{in} = 190 \text{ MeV}$  and  $T_{ex} = 180 \text{ MeV}$ .

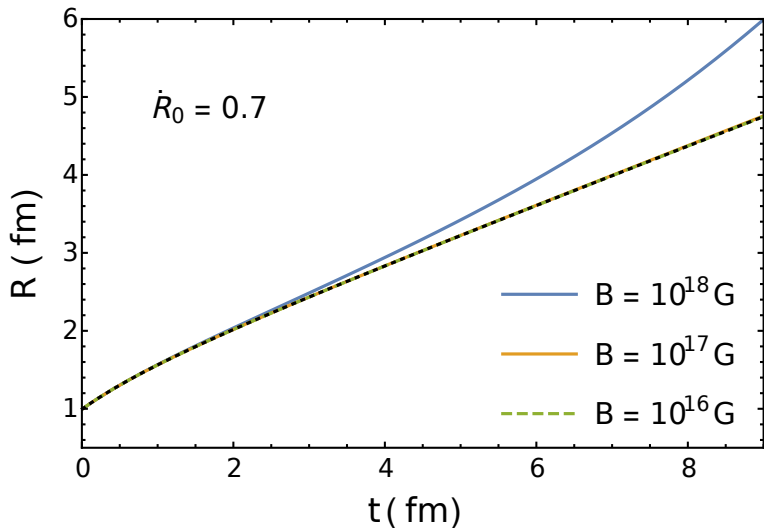


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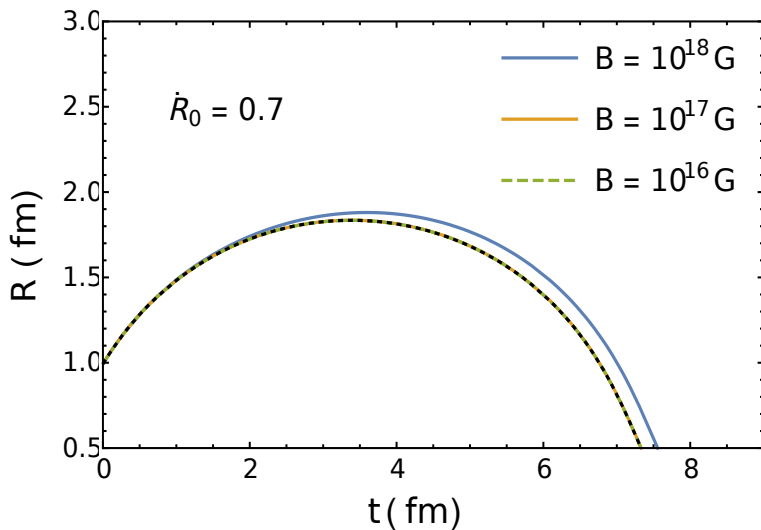


Figure: Collapse of hadron gas bubbles in QGP medium for  $T_{in} = 180$  MeV and  $T_{ex} = 190$  MeV.

# HG bubble collapse

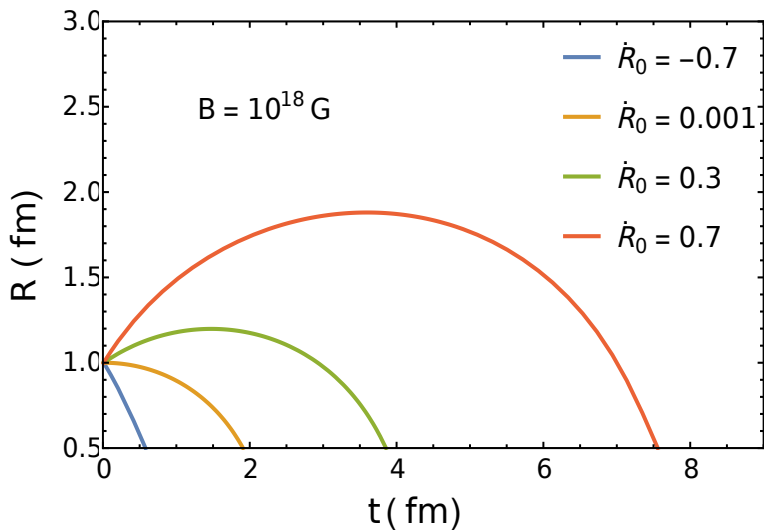


Figure: Collapse of hadron gas bubbles in QGP medium for  $T_{in} = 180$  MeV and  $T_{ex} = 190$  MeV.

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- Initial magnetic field increases the expansion rate of QGP bubbles
- Magnetic field at the initial stage of collision increases the collapse time of HG bubbles
- For large radial velocities, HG bubbles live longer

THANK YOU